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INFLUENCE OF IONIZING RADIATIONS AND ELECTRICAL  
OVERSTRESSINGS ON MOS (ME.) (U) RIT RESEARCH CORP  
ROCHESTER NY P S NEELAKANTASWAMY ET AL. JAN 87

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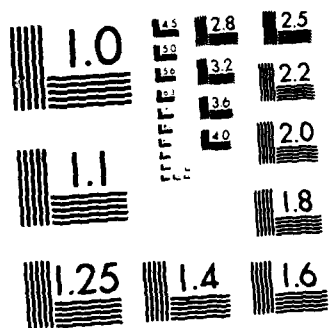
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REPORT DOCUMENTATION PAGE		READ INSTRUCTIONS BEFORE COMPLETING FORM
1. REPORT NUMBER RITRC-011, 11-11	2. GOVT ACCESSION NO.	3. RECIPIENT'S CATALOG NUMBER
4. TITLE (and Subtitle) Influence of Ionizing Radiations and Electrical Overstressings on MOS Devices: A Comparison		5. TYPE OF REPORT & PERIOD COVERED Tech. Report #11
7. AUTHOR(s) Perambur S. Neelakantaswamy Ibrahim R. Turkman		6. PERFORMING ORG. REPORT NUMBER
9. PERFORMING ORGANIZATION NAME AND ADDRESS RIT Research Corporation 75 Highpower Rd. Rochester, N.Y. 14623-3435		8. CONTRACT OR GRANT NUMBER(s) N00014-84-K-0532
11. CONTROLLING OFFICE NAME AND ADDRESS Office of Naval Research Arlington, VA. 22217		10. PROGRAM ELEMENT, PROJECT, TASK AREA & WORK UNIT NUMBERS NR 613-005
14. MONITORING AGENCY NAME & ADDRESS (if different from Controlling Office)		12. REPORT DATE January 1987
		13. NUMBER OF PAGES 23
		15. SECURITY CLASS. (of this report) Unclassified
		15a. DECLASSIFICATION/DOWNGRADING SCHEDULE
16. DISTRIBUTION STATEMENT (of this Report)		
Scientific Officer N00014 (1)		
Administrative Contract Officer S3305A (1)		
Director, Naval Research Laboratory N00173 (6)		
Defense Tech. Inform Center S47031 (12)		
17. DISTRIBUTION STATEMENT (of the abstract entered in Block 20, if different from Report)		
18. SUPPLEMENTARY NOTES		
19. KEY WORDS (Continue on reverse side if necessary and identify by block number) Ionizing Radiation Effects, MOS Devices, Electrical Overstressings, (EOS), Noise Characteristics.		
20. ABSTRACT (Continue on reverse side if necessary and identify by block number) Primary mode of failure and/or degradation of MOSFETs due to "oxide charge and surface-effects" can result either from ionizing radiations or from electrical overstressings. In either case, the resulting damage can be characterized by a global parametric degradation specified in terms of device noise characteristics. That is, the net effect of charge-trapping and the associated occupation of surface states can be viewed as random/fluctuation phenomena which manifest as the device noise. Thus a common noise model can be prescribed.		

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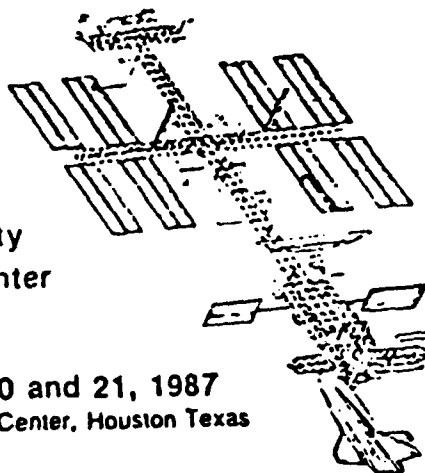
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INFLUENCE OF IONIZING RADIATIONS & ELECTRICAL  
OVERSTRESSINGS ON MOS DEVICES:  
A COMPARISON

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**ANALOGOUS INFLUENCE OF IONIZING RADIATIONS AND ELECTRICAL OVERSTRESSINGS:  
DAMAGE CHARACTERIZATION VIA NOISE PARAMETERS**

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**Abstract**

Primary mode of failure and/or degradation of MOSFETs due to 'oxide-charge and surface-effects' can result either from ionizing radiations or from electrical overstressings. In either case, the resulting damage can be characterized by a global parametric degradation specified in terms of device noise characteristics. That is, the net effect of charge-trapping and the associated occupation of surface states can be viewed as random/fluctuation phenomena which manifest as the device noise. Thus a common noise model can be prescribed to represent the analogous influence of ionizing radiations and electrical overstressings. Relevant theoretical results and measured data are presented.

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OVERSTRESSINGS: DAMAGE CHARACTERIZATION VIA NOISE PARAMETERS**

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**ABSTRACT**

Primary mode of failure and/or degradation of MOSFETs due to 'oxide-charge and surface-effects' can result either from ionizing radiations or from electrical overstressings. In either case, the resulting damage can be characterized by a global parametric degradation specified in terms of device noise characteristics. That is, the net effect of charge-trapping and the associated occupation of surface states can be viewed as random/fluctuation phenomena which manifest as the device noise. Thus a common noise model can be prescribed to represent the analogous influence of ionizing radiations and electrical overstressings. Relevant theoretical results and measured data are presented.

**INTRODUCTION**

The knowledge of common mechanisms involved in the degradation process(es) due to external stimuli, such as ionizing radiations and electrical overstressings, is useful, not only in understanding the interactive physics involved, but also will enable a common hardening technique (process/design) to achieve protection against these stimuli. Such studies will further indicate a one-to-one correlation (in quantifiable terms) between the intensity/magnitude of an ionizing influence and an electrical overstress which may cause the same extent of damage. This equivalence will enable substitution of test method(s) to simulate failure/degradation effects. 'Oxide-charge and surface-effects' [1,2] observed under the influence of ionizing radiations or electrical overstresses result from positive charge build-up in the gate-oxide due to radiation-induced (or EOS-induced) creation of electron-hole pairs; and the trapping of holes at the silicon-to-oxide interface alters the device parameters, namely, the transconductance ( $g_m$ ), MOS capacitance ( $C$ ) and the threshold voltage ( $V_T$ ). To understand the physics of these analogous effects observed, the mode(s) of energy transfer from the invasive external stimulus to the device interior, warrants unique modeling and analysis as discussed in this paper.

Inasmuch as all the degrading device parameters ( $g_m$ ,  $C$  and  $V_T$ ) are interdependent, the cohesive damage of the device would be assessed by an appropriate function which collectively represents the net physical damage due to external stimulus. It is presently demonstrated that noise characteristics can depict the global representation of stochastic variation [3] in charge-trapping and interface generation due to the external stimuli (ionizing radiations or EOS); relevant noise measurements of degraded devices can also be useful in accelerated test

procedures (using equivalent EOS to simulate ionizing radiations) adopted for life-time modeling strategies and in hardening effectiveness evaluation.

#### OXIDE-CHARGE & SURFACE EFFECTS

A MOS transistor can be looked at as a capacitor with the metal and semiconductor as the plates and the gate-oxide as the dielectric. Under ionizing radiation conditions, the ionization process is illustrated in Fig. 1. At  $t^- = 0$  (Fig. 1a), the condition prior to irradiation is shown. At  $t = 0$  (Fig. 1b), the ionizing energy is delivered to the oxide, and the electron-hole population is generated. Immediately after ionization, the process of electron-hole recombination will occur, but so will electron transport. But as electron mobility in the oxide at room temperature is approximately  $20 \text{ cm}^2/\text{V-sec}$ , and hole mobility is approximately  $2 \times 10^{-5} \text{ cm}^2/\text{V-sec}$ , under the applied voltage, any electrons that do not undergo recombination will be swept to the gate and removed in picoseconds, leaving behind the less mobile holes. These holes will begin a transport process toward the silicon-to-oxide interface as shown in Fig. 1e. Some holes will pass into Si, while others will become trapped at defect centers very near the interface of the gate oxide and the bulk silicon.

Fig. 2 depicts the shift in the C-V curve associated with the entire process and the resulting permanent shift due to the trapped charge buildup. In the case of the N-channel device, the trapped positive charge will continue to build up and, in effect, make it easier to create the N-channel (inversion layer). This will lower the threshold voltage (Fig. 3). The reversal of threshold shift is caused by the saturation of surface traps and interface state generation at the silicon-to-oxide boundary occurring at higher levels of ionizing radiations. This mechanism of interface state generation is not well understood at this time except for a simple theory that two different crystal structures (silicon and oxide) meet to form an interface having some irregularities, the number of which increases with increased irradiation. In the case of a corresponding P-channel device, the buildup will make it more difficult to create an inversion layer (in an enhancement mode P-channel transistor). The effect of ionizing radiation on a P-channel threshold is shown in Fig. 3. The net effects of ionizing radiation on a MOS device as a function of threshold shifts are therefore: N-channel devices are easier to turn on or can actually become depletion mode; and P-channel devices become more difficult to turn on.

Similar oxide-charge and surface-effects also appear when a MOS structure is subjected to an EOS, say by a (positive) high-voltage at the gate. During the high-voltage pulses, electrons are injected into the gate-oxide via Fowler-Nordheim tunneling from the Si substrate, and some fraction of the injected electrons then create electron-hole pairs in the bulk of the oxide through impact ionization (Fig. 4a). The resulting electrons and holes behave similarly to those generated by ionizing radiation in MOS structures under positive (worst-case) bias (Fig. 4b): Most of the electrons are swept out of the oxide while the



holes drift under the positive field toward the oxide-to-silicon interface where they may be removed or trapped. Some of these holes may also cause interface states to be produced. The resulting flatband-voltage shift and interface-state build-up can be depicted as in Fig. 2.

Ionizing radiations or electrical overstressings will also cause carrier mobility degradation because of the presence of trapped charges near the silicon-to-oxide interface and interface generation, of which interface generation is more dominant and it becomes negligible at lower levels of ionizing radiations/EOS. As the stressing levels are increased (about a million rad for ionizing radiations or 9MV/cm for EOS), mobility degradation will affect P and N-channel device performance, with increased interface states being the primary cause of degradation. This mobility degradation can be observed via transconductance ( $g_m$ ) measurements. Another performance problem induced by ionizing radiations or EOS is the increase in leakage current due to surface effects.

MOSFETs stressed by ionizing radiations or by EOS have the tendency to anneal. Annealing is the time-dependent detrapping of trapped charge at the silicon-to-oxide interface. It is sometimes referred to as a self-healing effect. However, the time constant involved is in the order of minutes to over one year, depending on the extent of damage, design-based on-chip protection and the type of processing. Though the surface states generated are relatively permanent, it can also be annealed with high temperatures ( $>125^\circ\text{C}$ ). Any lattice damage (interstitials, vacancies), however, is irreversible.

Experimental studies indicate that trapping of holes or oxide-silicon interface degeneration does not differ significantly between electrical overstressing and ionizing irradiations [4] despite the fact that holes are transported to the interface rather under high field conditions in EOS phenomenon; whereas, hole transport under ionizing irradiations is not field activated. Hence, it is evident that capture of hole by a trap at the interface is not a strong function of electric field in the oxide.

#### SINGLE-MODEL REPRESENTATION OF IONIZING RADIATION AND EOS EFFECTS

On the basis of aforesaid discussions, the identical effects observed in MOSFETs when subjected to ionizing radiations or EOS can be summarized as follows: 1) Shift in threshold voltage; 2) Change in oxide-capacitance; 3) Mobility ( $\mu$ ) degradation; 4) Change in transconductance; 5) Increase in leakage current; and 6) Annealing. These various parameters though can represent the degradation (either due to ionizing radiations or due to EOS) independently are, however, interdependent and explicitly related through analytical expressions. Therefore, it is possible to establish a general expression which uniquely represents the cohesive damage, irrespective of the nature of external stimulus. For this purpose the global effect of stochastic variations in charge-trapping and interface generation (under external stimulus) can be considered to model the net physical degradations observed. And as these stochastic fluctuations in the device-interior

manifest as the 'device-noise,' the desired modeling can be characterized by appropriate noise parameters of the device.

When a MOS device is subjected to external stress (either ionizing radiations or EOS), the corresponding induction of charge-trapping and generation of interface states can be equivalently represented by an input noise resistance  $R_N$  given by [5]

$$R_N \approx C_0 N_S (\mu_s / \mu_0)^2 \quad (1)$$

where  $C_0$  is a constant,  $N_S$  is the surface-state density and  $\mu_s / \mu_0$  refers to the field-effect mobility to low-field mobility ratio. Eqn. (1) indicates that  $R_N$  is directly proportional to  $N_S$  concurring with the experimental results due to Abovitz, et al [6]. Hence the time-dependent history of  $N_S$  as controlled by any external overstressings can be tracked via the assessment of  $R_N$ . Further, the field-effect mobility is also dependent on  $N_S$  and is therefore linked with  $g_m$  and  $V_T$  of the device. Explicitly,

$$\frac{\mu_s}{\mu_0} \approx \frac{1}{1 + \alpha N_S} \approx \frac{g_m}{g_{m0}} \approx \frac{1}{1 + \beta(V_G - V_T)} \quad (2)$$

Here,  $\alpha$  and  $\beta$  are constants and  $g_{m0}$  refers to the value of  $g_m$  under unstressed conditions and  $V_G$  denotes the applied gate potential.

From Eqns. (1) and (2), the following relation can be obtained:

$$\frac{\Delta R_N}{R_N} \approx \frac{\Delta g_m}{g_m} \left[ 2 - \frac{1}{\frac{1 - \Delta V_T}{V_T} \frac{1}{(V_G - V_T)\beta}} \right] \quad (3)$$

The constant  $\beta$  has the approximate values of 0.138 and 0.308 for the N-channel and P-channel MOSFETs, respectively [4]. More generally,  $V_G$  can be expressed in terms of the electric field across the gate-oxide, namely,  $E_G$ . That is,  $V_G = E_G t_{ox}$ , where  $t_{ox}$  is the gate-oxide thickness. While  $E_G$  refers to the electric field intensity corresponding to an electrical overstress (EOS) phenomenon, it is possible to establish an equivalent  $E_G$  to represent the ionizing radiation dosage, which produces the same extent of degradation expressed via noise parameter of eqn. (3). Let  $D_I$  be the dosage delivered (or, absorbed dosage) to an oxide-gate, which through ionization process creates a hole density (area density) of  $Q_R$  equal to  $K_I t_{ox} D_I F(E_G)$  where  $K_I$  is the infinite field ionization coefficient [4]

equal to  $1.22 \times 10^{-6} \text{ C CM}^{-3} \text{ rad}^{-1}$  ( $\text{SiO}_2$ ) and  $F(E_G)$  is the E-field dependent charge-yield parameter [4] with approximate value of 0.83 at  $E_G = 1 \text{ MV/cm}$ . If the same hole-density of  $Q_R$  has to be stimulated by an electrical overstress phenomenon (via high field injection of electron current density through Fowler-Nordheim tunneling by the gate-oxide field-intensity,  $E_G$ ), the corresponding current density ( $j$ ) can be expressed (by neglecting space-charge effects) as equal to  $A E_G^2 \exp(-B/E_G)$  where A and B are constants. The best estimates of [4] A and B are  $2 \times 10^6 \text{ amperes/(MV)}^2$  and  $238 \text{ MV/cm}$ , respectively. Therefore EOS-equivalent of  $Q_R$  can be written as

$$Q_R (\text{EOS}) = j \alpha t_{ox} \Delta t \quad (4)$$

where  $\alpha$  is the probability per unit length that an injected electron will create an electron-hole pair and is equal to  $\alpha_0 \exp(-H/E_G)$  where  $\alpha_0 = 6.5 \times 10^{11} \text{ cm}^{-1}$  and  $H \sim 180 \text{ MV/cm}$  [4].

Further,  $\Delta t$  in eqn. (4) specifies the duration of EOS event. Assuming that lightning function of the form  $e_1(t) = E_G [\exp(-Ct) - \exp(-Dt)]$  to represent the transient electrical overstressing, the duration,  $\Delta t$  is given by

$$\Delta t = [(D-C)/CD] [\exp(-Ct_m) - \exp(-Dt_m)]^{-1} \quad (5)$$

where  $t_m$  is the rise-time of the transient equal to  $\ln(D/C)/(D-C)$ . The values of C and D can be explicitly specified for a given type of EOS event, such as human-body ESD model, etc., and are dependent on the peak value of the stressing potential,  $V_G$ .

Combining eqn. (4) and eqn. (5), the equivalent dosage  $D_I$  can be expressed as

$$D_I = \alpha_0 A E_G^2 [(D-C)/CD] [\exp(-Ct_m) - \exp(-Dt_m)]^{-1} \exp[-(H+B)/E_G] \quad (6)$$

Hence, using approximate values for the parameters discussed previously, the expression for  $D_I$  reduces to

$$D_I = 10^{24} (V_G/t_{ox})^2 \exp(-418 t_{ox}/V_G) \Delta t (V_G) \quad (7)$$

where  $V_G$  is expressed in MV,  $t_{ox}$  in cm and  $\Delta t$  is functionally dependent on  $V_G$ .

Considering a typical electrical overstressing due to an electrostatic discharge (ESD) of subcatastrophic level (say  $V_G = 50V$  peak) from a finger tip (Fig. 5) across a gate-oxide of thickness 30 nm over a pulse-duration,  $\Delta t = 10n$  sec, the corresponding equivalent radiation dosage is approximately equal to 36 M rad ( $SiO_2$ ). That is, this ESD event would introduce as many holes into the oxide as a 36 M rad ( $SiO_2$ ) of ionizing radiation.

#### RELATIVE NOISE PERFORMANCE UNDER EOS AND IONIZING IRRADIATIONS

For a given injected electron fluence ( $Q_T = j\Delta t$ ), the relative damage introduced in the MOSFET by an EOS and an ionizing radiation can be estimated as follows: By virtue of one-to-one equivalence between the magnitude of EOS and ionizing dosage, the relative damage expressed, say, in terms of threshold shift  $\Delta V_T/V_T$  can be written as a linear proportionality relation of the form  $\Delta V_T(EOS)/V_T = K \Delta V_T(RAD)/V_T$  where  $K$  is a constant.

Considering the results due to Boesch and McGarrity [4], for a given amount of injected electron fluence ( $1.9 \times 10^{-5} C/cm^2$ ), the theoretical and experimental results corresponding to high field stressing of 9MV/cm) and  $^{60}Co$  irradiation ( $10^4$  rad  $SiO_2$ ) on a  $t_{ox} = 1000$  Å MOS structure, the value of  $K$  is found to be approximately equal to 2.3; or, in general,  $K > 1$ . Hence, using the linear relation between  $\Delta V_T(EOS)/V_T$  and  $\Delta V_T(RAD)/V_T$ , it can be shown that  $\Delta R_N(EOS)/R_N$  is nearly equal to  $(4/3)\Delta R_N(RAD)/R_N$ . In other words the damage, manifesting as the device noise under EOS injecting a given amount of electron fluence into the gate, is approximately 25% more, for the same extent of electron fluence injected by a radiation source. Typical noise parameter ( $\rho = \frac{\Delta R_N/R_N}{\Delta g_m/g_m}$ ) variations as functions of radiation dosage for a P and N channel MOSFETs are shown in Figs. 6 and 7. Corresponding variations of threshold voltage,  $V_T$  are also depicted in Figs. 6 and 7, from which it can be observed that the noise parameter follows the trend of  $V_T$  variations(s).

Will radiation hardening concurrently improve static-protection or vice versa? The observed similarities suggest the possibility of formulating a one-to-one equivalence of modeling of radiation damage versus EOS effects from which it can be extrapolated that any scheme that is implemented (either via processing or via design methods) to prevent/reduce radiation-induced (deleterious) effects may also subdue the influences of EOS effects. In other words radiation hardening schemes (process/design) with a few optimization changes may provide dual protection to prevent/reduce gate-oxide damages arising from ionizing radiations or from EOS. In order to achieve effective dual protection through optimization of process/design techniques, basic

research is required to determine this ionization radiation-to-EOS equivalence so that the common-to-both type of damages(s) in the gate-oxide can be effectively prevented through optimization procedures.

#### CONCLUSIONS

This work provides a basic insight into the problem of a comparative study relating ionizing radiations and EOS effects on MOS devices. The results indicate a strong correlation between the two effects cited, which suggests the feasibility of designing common countermeasures, as well as adopting substitutions in the analysis and/or simulation techniques.

#### ACKNOWLEDGEMENT

This work was supported by a grant from the Office of Naval Research (No. 613-005) which is gratefully acknowledged.

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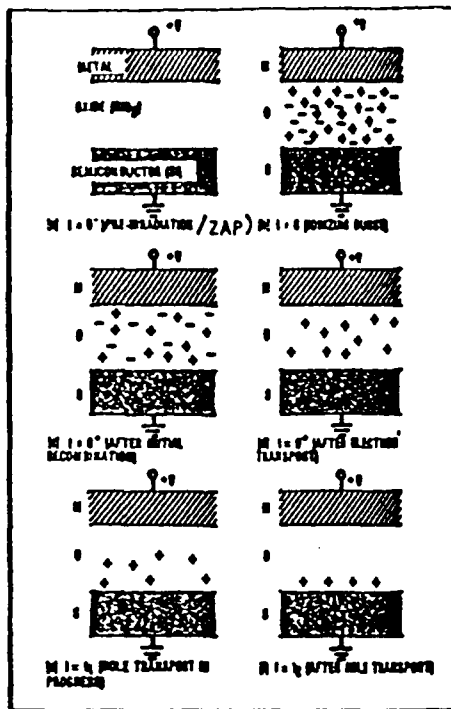


FIG. 1 CARRIER TRANSPORT MECHANISM

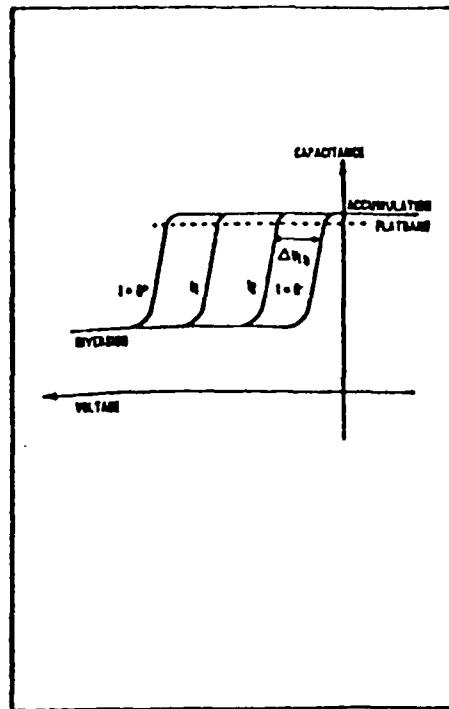


FIG. 2 C-V CURVES FOR CONDITIONS IN Fig. 1

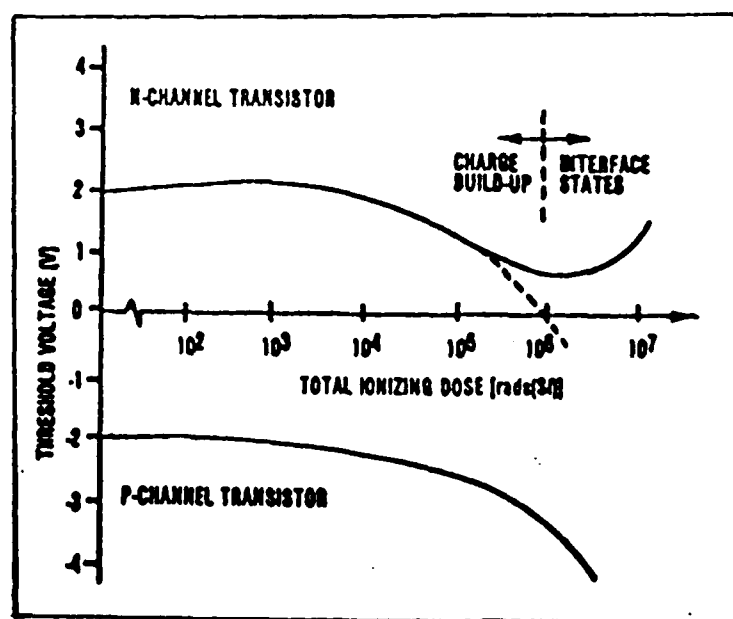


FIG. 3 RADIATION EFFECTS ON P AND N CHANNEL MOSFETS.

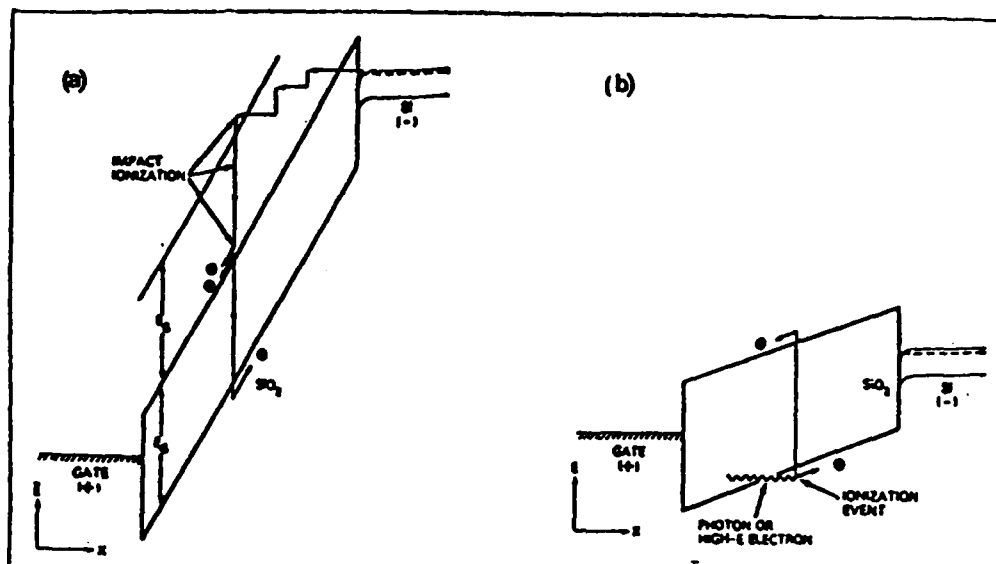


FIG. 4 CREATION OF ELECTRON - HOLE PAIRS DUE TO: (A) HIGH-FIELD INJECTION AND (B) IONIZING RADIATION

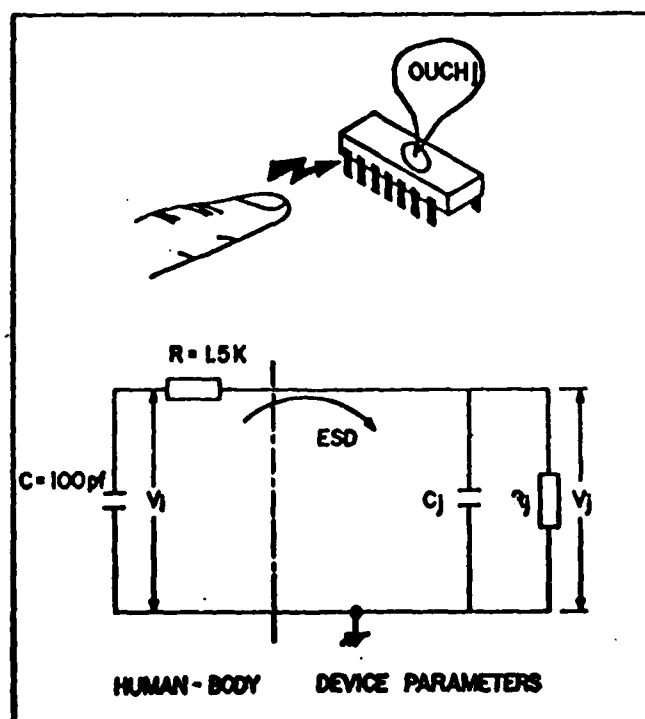


FIG. 5 ESD: HUMAN-BODY MODEL

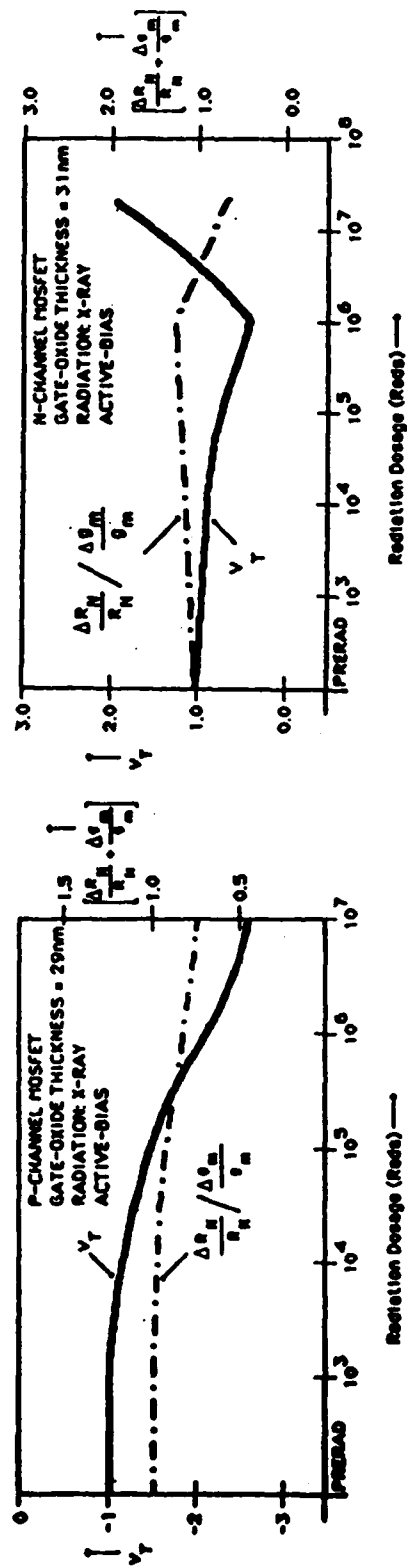


FIG. 6 RADIATION DOSAGE VERSUS THRESHOLD VOLTAGE ( $V_T$ ) AND NOISE PARAMETER ( $\Delta R_N/R_N + \Delta g_m/g_m$ )

FIG. 7 RADIATION DOSAGE VERSUS THRESHOLD VOLTAGE ( $V_T$ ) AND NOISE PARAMETER ( $\Delta R_N/R_N + \Delta g_m/g_m$ )



## NOTE

### NOISE CHARACTERISTICS OF IONIZING-RADIATION STRESSED MOSFET DEVICES

#### INTRODUCTION

It has been known for years that ionizing radiations can change the electrical properties of solid state devices, leading to possible system failure [1]. In particular, gamma rays, X-rays and neutron bombardment have proven most harmful. Among the LSI devices, MOS circuits are highly sensitive to damages under critical radiation environments. The primary failure mode and/or degradation of MOSFETs resulting from ionizing radiations is due to the 'oxide-charge and surface-effects,' [2], occurring in the gate-oxide and/or field-oxide regions. The effects of ionizing radiations are mainly threshold voltage shift and channel mobility degradation caused by the creation of electron-hole pairs and trapping of holes at the Si-SiO<sub>2</sub> interface [2]. The net effect of charge-trapping and the associated occupation of surface states can also be viewed as a random/stochastic phenomenon which can be characterized by a global noise parameter. Such a representation/model will be useful to study the noise performance of the device under radiation environments as indicated in the present work.

#### NOISE-MODEL

The global effect of 'oxide-charge and surface-effects' described above can manifest as the device noise which can be quantitatively represented by a noise-model as described below:

Following the analysis by Leventhal [3], the effect of charge-trap induction and generation of interface states can be equivalently represented by an input noise resistance  $R_N$  given by

$$R_N = C \left( \frac{\mu_S}{\mu_0} \right)^2 N_S \quad (1)$$

where  $C$  is a constant of proportionality and  $N_S$  is the surface-state density;  $\mu_S/\mu_0$  refers to the field-effect mobility to low field-effect mobility ratio. Further, inasmuch as field-effect mobility is also dependent on  $N_S$  and is, therefore, linked with the device transconductance ( $g_m$ ) and the threshold voltage ( $V_T$ ) [4,5], the following relation can be obtained [6] from eqn. 1. (Note the typographical errors in [6]: eqns. 1 and 3 of [6] should read as eqns. 1 and 2 of the present paper, respectively.)

$$\frac{\Delta R_N}{R_N} = \frac{\Delta g_m}{g_m} \left[ 2 - \frac{1}{1 - \frac{\Delta V_T}{V_T}} \cdot \frac{1}{(V_G - V_T)\beta} \right] \quad (2)$$

where  $\Delta V_T$  is the threshold-voltage shift and  $\beta$  is a constant, approximately equal to 0.138 and 0.308 for the N-channel and P-channel MOSFETs, respectively [5].

The quantity  $V_G$  in eqn. 2 represents an 'equivalent gate-potential' which would inject the same electron-fluence into the gate equal to that injected by the ionization irradiation.  $V_G$  can also be expressed in terms

of an 'equivalent electric field' across the gate-oxide, namely  $E_G$ . That is,  $V_G = E_G t_{ox}$ , where  $t_{ox}$  is the gate-oxide thickness.

This equivalent electric stress parameter  $E_G$ , can be specified explicitly in terms of the radiation dosage,  $D$ , assuming that, over a duration of  $\Delta t$ , the effect of  $E_G$  or  $D$  is to inject the same extent of electron fluence, namely  $j\Delta t$ , where  $j$  is the current density. Hence, relevant analysis yields,

$$D \approx A E_G^2 \Delta t \exp (-B/E_G) \quad (3)$$

where  $A$  and  $B$  are constants approximately equal to  $10^{24}$  and 418, respectively [7], if  $E_G$  is expressed in MV/cm. Thus for a given dosage level of  $D$ , using eqns. 2 and 3, the noise performance of the device can be decided quantitatively.

#### TEST STUDIES & CONCLUSIONS

Variations of the noise parameter  $\frac{\Delta R_N}{R_N} / \frac{\Delta g_m}{g_m}$  corresponding to a P-

channel and an N-channel MOSFETs as functions of the radiation dosage (X-ray), are depicted in Figs. 1 and 2. Also shown in Figs. 1 and 2 are the threshold voltage shifts in the P- and N-channel MOSFETs [8]. The results on noise parameter presented in Figs. 1 and 2 are calculated via eqns. 2 and 3, using the available data on threshold voltage shifts versus radiation dosage [8]. From the results shown, the following can be inferred:

1. Damage introduced by ionizing radiations in a semiconductor device (such as MOSFET) can also be characterized by the noise performance of the device.
2. The noise parameter, as a function of radiation dosage tends to track closely the variation of  $V_T$  with respect to the dosage level. This is true for both P- and N-channel MOSFETs (Figs. 1 & 2). The percentage shift in the magnitude of noise parameter, for a given level of radiation dosage is, however, less than the corresponding percentage shift in the threshold voltage. Referring to Fig. 1, (P-MOS) for a dosage level of  $10^7$  rad (Si),  $V_T$  shifts by 170%, whereas the noise parameter changes only by 26%. Similarly in Fig. 2 (N-MOS),  $V_T$  shifts for  $10^6$  and  $10^7$  rad (Si) are -52% and +42%, respectively. However, the corresponding noise parameter shifts are +28% and -8%, respectively. Thus, for a given level of irradiation, the variation in  $V_T$  is more overwhelming than changes in noise performance.
3. Nevertheless, in low noise applications of the device, the influence of ionizing radiations should be duly accounted for in the system design as noise performance degradation is inevitable as a result of radiation-induced oxide effects. While 'single event' upsets due to ionizing radiation usually cause concern in digital circuits, noise performance degradation due to cumulative/total ionizing irradiations may require specific attention in linear devices. Especially, as the device is stressed repeatedly, the device damage (noise performance degradation) will cumulatively increase. Such endochronic degradation response would be detrimental for low-noise system operation. The present analysis is useful in the relevant studies.

4. The simple model presented here provides a quantitative approach to determine the noise performance of a MOSFET ionizing irradiations. The relevant calculations are useful to determine the extent of radiation hardening required to achieve a given level of low noise performance of the device under ionizing radiation environments. And, noise monitoring can serve as an adjunct support to conventional  $V_T$  and  $C_{ox}$  estimations adopted in hardness assurance efforts.
5. The present work models only the effect(s) of ionizing radiations on the device-noise. Should the geometrical parameters (such as the channel length) change, the transconductance would be significantly affected (especially in short-channel devices) and the relevant noise-model will be more involved. Related studies are in progress.
6. It can be shown that  $\Delta g_m/g_m = \Delta V_T/(V_G - V_T)$ . Hence, the noise-parameter profiles of Figs. 1 and 2 will remain the same (except for a scale-factor) if the noise-parameter is normalized with respect to threshold-voltage shift.

Acknowledgement: This work is supported by a grant by the Office of Naval Research (No. 613-005) which is gratefully acknowledged.

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**CAPTIONS FOR THE DIAGRAMS**

**Fig. 1** Threshold voltage and noise parameter versus total radiation dosage (P-Channel MOSFET).

**Fig. 2** Threshold voltage and noise parameter versus total radiation dosage (N-Channel MOSFET).

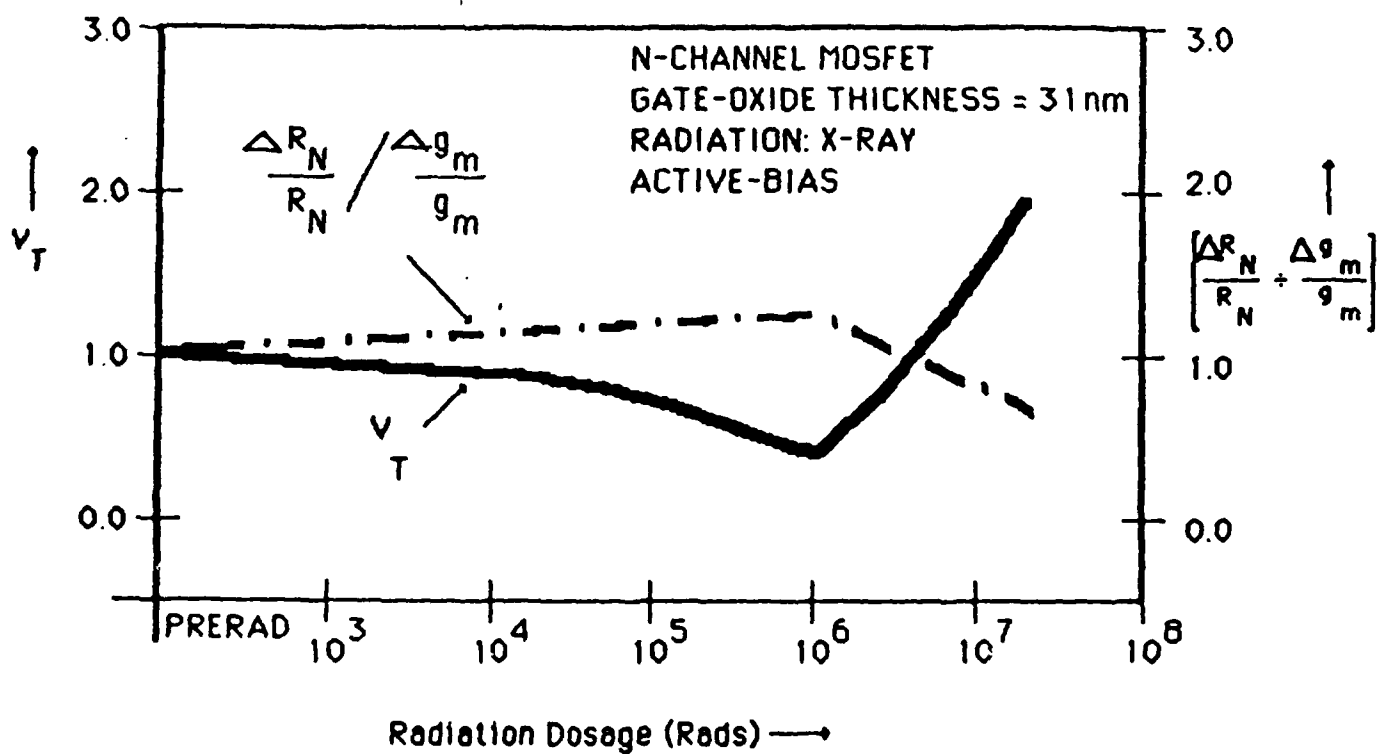


FIG. 2 THRESHOLD VOLTAGE ( $V_T$ ) AND NOISE PARAMETER ( $\Delta R_N / R_N + \Delta g_m / g_m$ )  
VERSUS RADIATION DOSAGE  
N-CHANNEL MOSFET



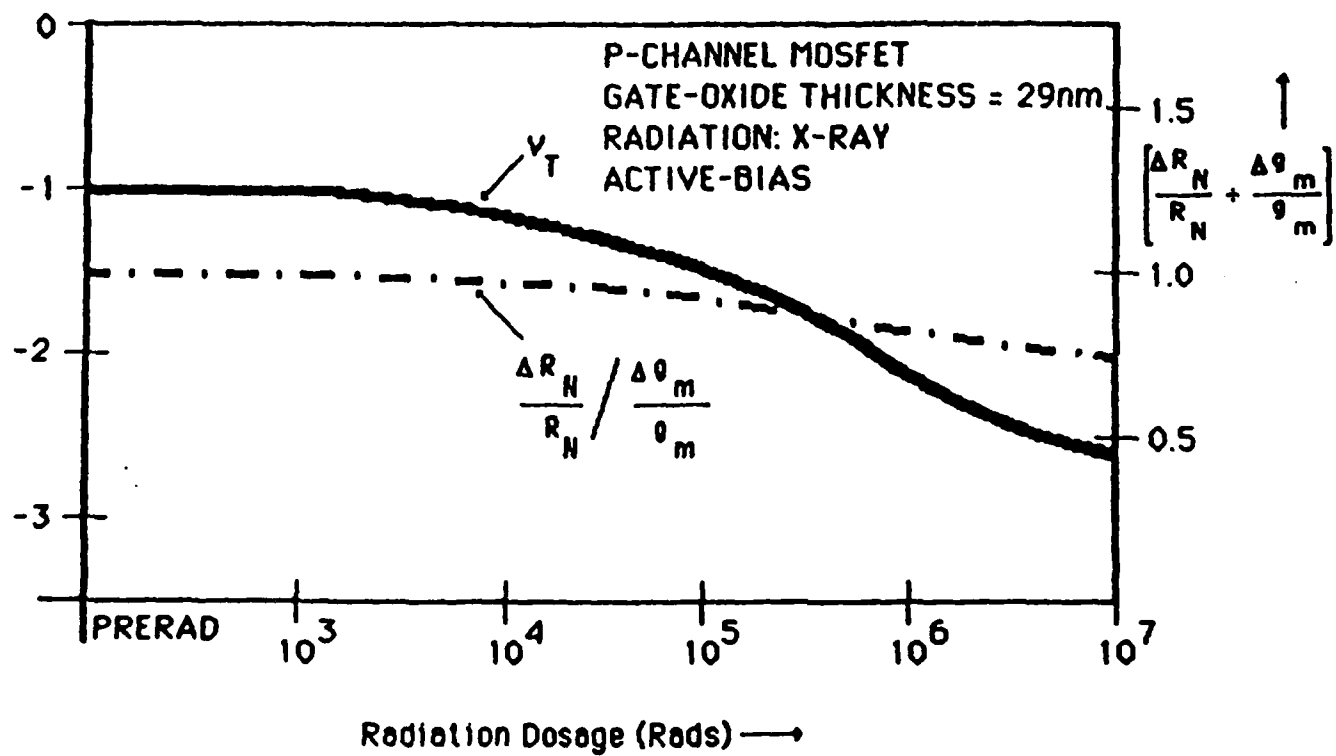


FIG. 1 THRESHOLD VOLTAGE ( $V_T$ ) AND NOISE PARAMETER ( $\Delta R_N / R_N \div \Delta g_m / g_m$ )  
VERSUS RADIATION DOSAGE  
P-CHANNEL MOSFET

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